


Conservation laws and shock waves.

$$\begin{cases} \partial_t \rho + \nabla \cdot (\rho u) = 0, \\ \partial_t (\rho u) + \nabla \cdot (\rho u \otimes u + pI) = 0, \\ \partial_t (\rho E) + \nabla \cdot (\rho E u + p u) = 0, \quad p = p(\rho, e), \quad \rho E = \frac{1}{2} \rho |u|^2 + \rho e. \end{cases}$$

← Euler: $\partial_t v + \nabla \cdot (f(v)) = 0$, → conservation law. (hyperbolic).

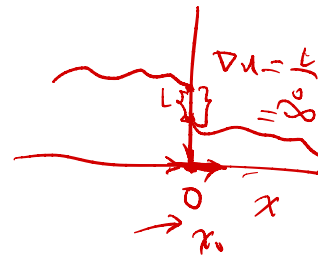
$v = \begin{pmatrix} \rho \\ \rho u \\ \rho E \end{pmatrix}$, $(\rho u, \rho u \otimes u + pI, \rho E u + p u) = f(v)$

NS: $\partial_t v + \nabla \cdot (f(v)) = \nabla \cdot (B(u, \varepsilon) \nabla u)$. $\varepsilon \rightarrow 0 \rightarrow (2n+1)$. $n+2$. $n+2$. toy model → 

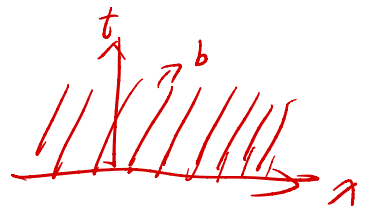
Schrodinger / incompress Euler / NS ↔ compressible Euler / NS

↓ shock

- $\Delta |\nabla u| = \infty$
- $x = x_0 = f(x)$
- RH condition
- entropy condition



method: characteristic curve



$u_t + b u_x = 0$

transport equation

Toy model:

$\partial_t u + u \cdot \nabla u + \nabla p = 0$ - Incompressible Euler

$\|u\|_{H^s} \rightarrow \infty \Leftrightarrow \int_0^{T^*} \|w\|_{L^\infty} = \infty$, $w = \nabla \times u$
 $t \rightarrow T^*$, $L^1 L^\infty(0, T^*)$

Evans chapter 3.

conservation law

$\partial_t w + u \cdot \nabla w = w \cdot \nabla u$ - vorticity equation

→ axisymmetric



$$\rightarrow \left[T = \frac{c}{H} \right]$$

1D toy model of vorticity equation.

$$\partial_t w = w H w$$

$$w = w_0$$

$$w = \frac{w_0}{(1 + t H w_0)^2 + t^2 w_0^2}$$

$$t \rightarrow \frac{1}{H w_0}$$

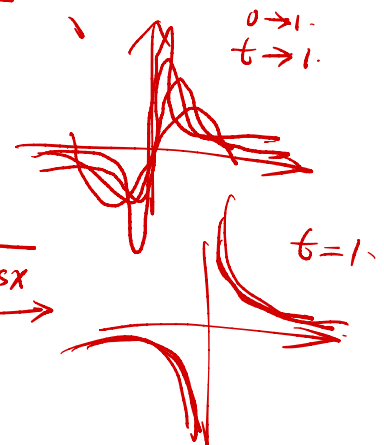
$$w_0 = \sin x$$

$$w = \frac{\sin x}{1 + t^2 - 2t \cos x}$$

$$t \rightarrow 0 \rightarrow 1$$

$$\sin x$$

$$2 - 2 \cos x \rightarrow x=0 \quad (|w(t)|)_{t \rightarrow \infty} = \infty$$



Elginder 2021 $c^{1,\alpha} \rightarrow$ blow up

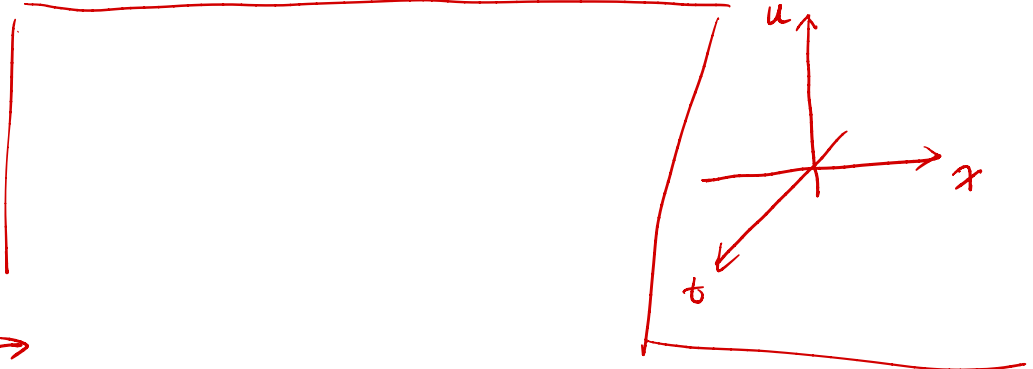
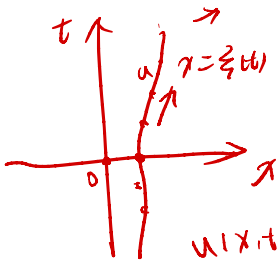
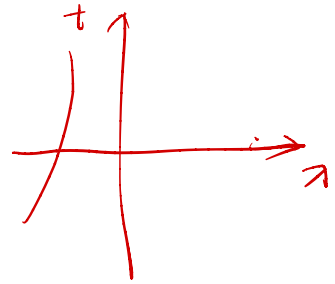


1d conservation law

$$\partial_t u + \nabla \cdot (f(u))$$

$$\partial_t u + \partial_x (f(u)) = 0$$

$$\Rightarrow u_t + (f(u))_x = 0$$



$$\textcircled{1} \quad x = \xi(t)$$

$$\frac{d}{dt} u(\xi(t), t) \equiv 0$$

$$\xi'(t) u_x(\xi(t), t) + u_t(\xi(t), t) = 0$$

$u \in C$

$$\partial_t u + f'(u) u_x$$

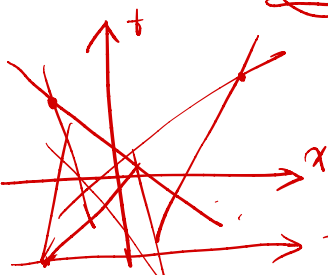
$$(u_t + \xi' u_x)(\xi(t), t) = 0$$

$$\xi' = f'(u) = \lambda(u)$$

$\rightarrow x = \xi(t)$ is a line

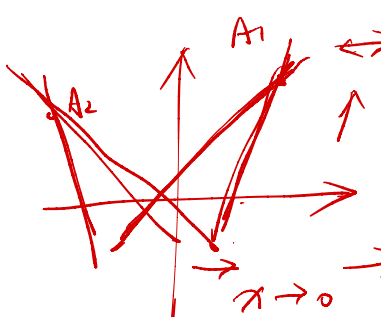
$$\xi'(t) = f'(u)(\xi(t), t) = \lambda(u)(\xi(t), t) \equiv \lambda(u)(\xi(0), 0)$$

u



piecewise smooth solution
(in distributional sense)





$x \rightarrow 0$ (compression wave) shock: $f''(u) \geq 0$.

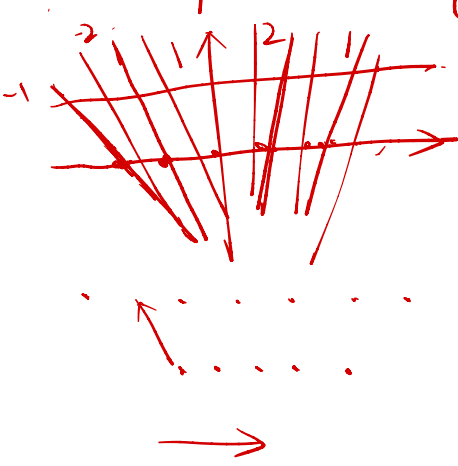
$A_1: 1 \rightarrow 2$

$$\left(\frac{\xi}{x}\right)' \geq 0$$

$A_2: 2 \rightarrow 1$

(expansion wave) (rarefaction wave formation)

$$\xi' = f'(u) \rightarrow f''(u) \left(\frac{\xi}{x}\right)' \geq 0$$

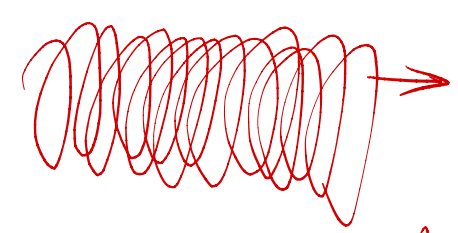


$$f''(u) < 0$$

$$f''(u) \leq 0$$

$$f''(u) \geq 0$$

f is convex function



Concave

compression

expansion

$$f''(u) < 0$$

$$f''(u) > 0$$

$$f(u) = \frac{1}{3}u^3$$

$$f'(u) = u$$

$$u > 0$$

$$u = 0$$

$$u < 0$$



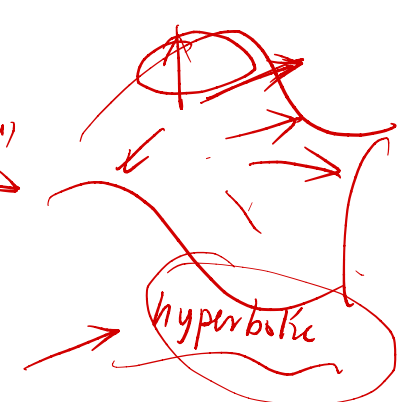
convex concave
 $f''(u) > 0$ $f''(u) < 0$

conservation system

$$f''(u) > 0$$

$$f''(u) < 0$$

$$f''(u)$$



Hopf equation:
 (inviscid equation)
 Burgers

critical
 $u_t + u u_x = 0$

$$f(u) = \frac{1}{2} u^2$$

$$f''(u) = 1 \rightarrow \text{convex}$$

$$(u_t)_t + \nabla \cdot (u u_x)$$

$$u_t + \nabla u^2 = 0$$



Some examples of shock waves and rarefaction waves:

RH condition: \rightarrow motion of shock wave.

entropy condition \rightarrow uniqueness solution.

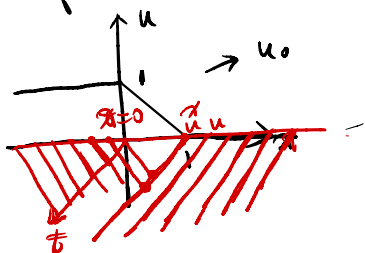
(inviscid Burgers)

We consider Hopf equation:

$$\begin{cases} u_t + uu_x = 0 \\ u(0) = u_0 \end{cases}$$

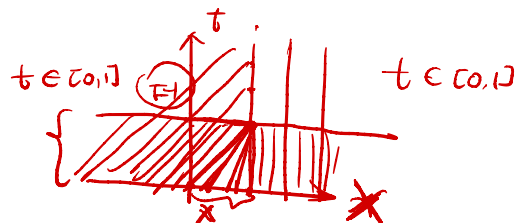
$$u_0 = \begin{cases} 1 & x \leq 0 \\ 1-x & x \in (0, 1] \\ 0 & x \geq 1 \end{cases}$$

$u_0 \geq 0$



$$\xi'(t) = 1-x \quad (u)' = u = 1$$

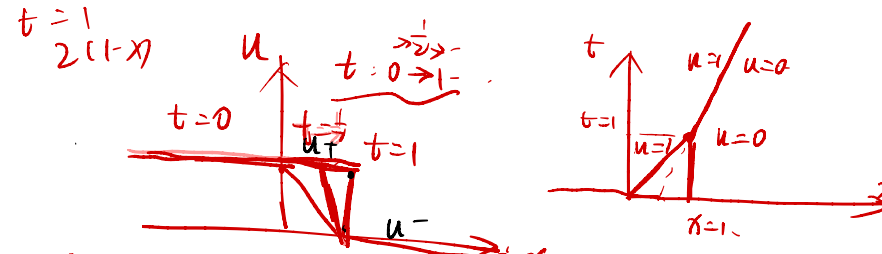
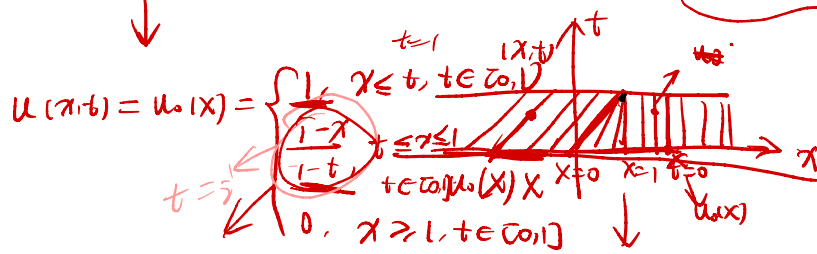
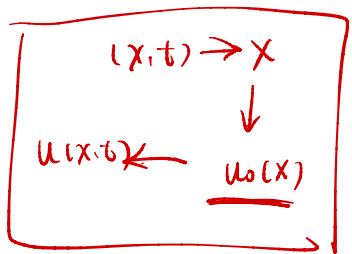
$$\xi'(t) = f'(u) = u = 0/1$$



$$x = \begin{cases} x+t & x \leq 0 \\ x+(1-x)t & x \in (0, 1] \\ x & x > 1 \end{cases}$$

$(X, t) \rightarrow x$

$$X = \begin{cases} x-t & x \leq t, t \in (0, 1] \\ \frac{x-t}{1-t} & t \in (x, 1+t) \\ x & x > 1, t \in (0, 1] \end{cases}$$



② Rankine-Hugoniot condition:
RH

$$\xi'(t) = \frac{f(u^+) - f(u^-)}{u^+ - u^-} = \frac{\frac{1}{2}u^2 - \frac{1}{2}u^2}{u^+ - u^-} = \frac{u^+ + u^-}{2} = \frac{1}{2}$$

$f(u) = \frac{1}{2}u^2$

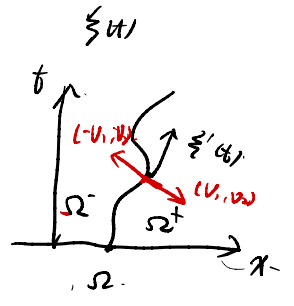
Theorem: suppose (u^+, u^-) , (jump condition), u satisfies conservation law: $u_t + f(u)_x = 0$,

then $(u^+, u^-) \rightarrow \xi'(t)$ satisfies: $\xi'(t) = \frac{f(u^+) - f(u^-)}{u^+ - u^-}$

Proof:

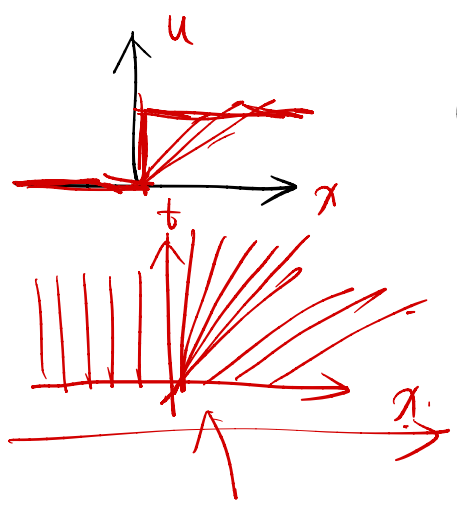
$$\begin{aligned}
 0 &= \int_0^{+\infty} \int_{-\infty}^{+\infty} (u u_x + f(u) v_x) dx dt \\
 &= \int_{\Omega^-} + \int_{\Omega^+} (u u_x + f(u) v_x) dx dt \\
 &= \int_{\Omega^-} + \int_{\Omega^+} (f(u), u) \cdot \nabla_{x,t} v dx dt \\
 &= \int_{\Omega^-} + \int_{\Omega^+} \underbrace{\nabla_{x,i} (f(u), u)}_{(u_t + f(u)_x)} v dx dt + \int_{\partial \Omega^-} v (f(u), u) \cdot n dx dt \\
 &\quad + \int_{\partial \Omega^+} v (f(u), u) \cdot n dx dt \\
 &= \int_{\mathbb{R}} v (f(u^+), u^+) \cdot (v_1, v_2) + \int v (f(u^-), u^-) \cdot (-v_1, -v_2) \\
 &= \int_{\mathbb{R}} v (f(u^+) - f(u^-), u^+ - u^-) \cdot (v_1, v_2)
 \end{aligned}$$

$v \in C_c^\infty(\mathbb{R} \times [0, \infty))$
 $v(x, 0) = 0$

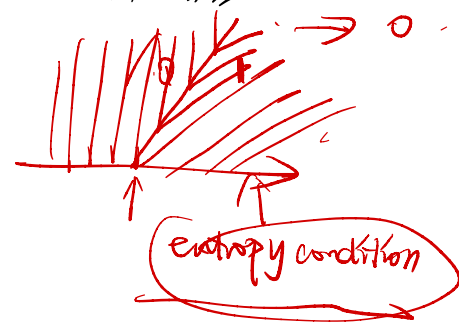


$$\frac{f(u^+) - f(u^-)}{u^+ - u^-} = -\frac{v_2}{v_1} = \xi'(t)$$

②



$$u_0(x) = \begin{cases} 0, & x \leq 0 \\ 1, & x \geq 0 \end{cases}$$



$$\begin{array}{ccc}
 & \rightarrow & \\
 u^+ & > \xi'(t) & > u^- \\
 \hline
 0 & & 1
 \end{array}$$